

Lecture 2

Thursday, September 17, 2020 9:58 AM

Recall:

Integrable system \Leftrightarrow compatibility conditions of overdetermined system of linear problems

Vast majority of examples

Ex Lax equation

$$\dot{L} = [A, L] \quad L(t) \text{ linear operator}$$

KdV a) $u_t + \frac{3}{2}uu_x - \frac{1}{4}u_{xxx} \quad L = -\partial_x^2 + u(x,t)$

NLS b) $ir_t = r_{xx} \pm |r|^2 r \quad L = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \partial_x + \begin{pmatrix} 0 & r \\ \pm \bar{r} & 0 \end{pmatrix}$

Toda c) $\ddot{x}_n = e^{x_{n-1} - x_n} - e^{x_n - x_{n+1}}$

$$(L\psi)_n = \psi_{n+1} + v_n \psi_n + \psi_{n-1}$$

$$\psi = (\psi_n)$$

Zero-curvature equations

$$0 = [\partial_z - U, \partial_y - V]$$

$$U = \begin{pmatrix} v & 1 \\ \lambda^{-1} & -v \end{pmatrix}$$

$$V = \begin{pmatrix} 0 & \lambda e^u \\ e^{-u} & 0 \end{pmatrix}$$

$$2u_{z\bar{y}} = e^u - e^{-u} \quad u \rightarrow iu$$

d) sine-gordon

$$u_{z\bar{y}} = \sin u$$

e) (L, A, B) triple

$$[L, \partial_t - A] = BL$$

... ..

$$[L, \partial_t - A] = BL$$

$$L = \partial \bar{\partial} + u, \quad A = \partial^3 + v \partial$$

$$\partial = \partial_z \quad \bar{\partial} = \partial_{\bar{z}}$$

$$\dot{L} = [A, L] + BL$$

$$[\partial^3 + v \partial, \partial \bar{\partial} + u] = 3u_z \partial^2 + 3u_{z\bar{z}} \partial + u_{z\bar{z}z}$$

$$- v_{\bar{z}} \partial^2 - v_z \partial \bar{\partial} - v_{z\bar{z}} \partial + v u_z$$

$$\dot{u} = 3u_z \partial^2 + 3u_{z\bar{z}} \partial + u_{z\bar{z}z} - v_{\bar{z}} \partial^2 - v_z \partial \bar{\partial} - v_{z\bar{z}} \partial + v u_z$$

$$B(\partial \bar{\partial} + u) = u_{z\bar{z}z} + v u_z + v_z u$$

$$3u_z = v_{\bar{z}}$$

$$B = v_z$$

$$= u_{z\bar{z}z} + (v u)_z$$

Equation Novikov-Veselov

$$L\psi = 0$$

$$(\partial_t - A)\psi = 0$$

$$(L - [AL])\psi = 0$$

=

If there are enough $\psi \Rightarrow$ Lax equation

If the ideal of operators annihilating ψ is generated by $L \Rightarrow$ L, A, B triple

!!! Linear operators depend on a "spectral parameter"
(explicit or latent form)

\mathcal{E}_x

Toda lattice

$$\ddot{x}_n = e^{x_{n-1} - x_n} - e^{x_n - x_{n+1}}$$

$$x_{n+N} = x_n + l$$

N periodic Toda lattice

$$(L\psi)_n = \psi_{n+1} + v_n \psi_n + c_n \psi_{n-1}$$

$$v_n = v_{n+N}$$

$$c_n = e^{x_n - x_{n-1}} = c_{n+N}$$

$$\psi_{n+1} = v \psi_n + c_n \psi_{n-1} \quad c_n = c_{n+N}$$

$$c_n = e^{\lambda_n - \lambda_{n-1}} = c_{n+N}$$

$$\psi_{n+N} = w \psi_n$$

$$\begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \\ \psi_N \end{pmatrix}$$

$$\psi_{N+1} = v \psi_1 \dots$$

$$\psi_0 = w^{-1} \psi_N$$

$$\begin{pmatrix} \vdots \\ \psi_0 \\ \psi_1 \\ \vdots \\ \psi_N \end{pmatrix} \quad L(w) = \begin{pmatrix} w_0 & \dots & w_N \\ \vdots & \ddots & \vdots \\ c_n v_n^{-1} & \dots & \dots \\ \vdots & \dots & \vdots \\ c_N \end{pmatrix}$$

Next goal is to define and then "solve"

- Phase space



space
of operators

$[A, L]$ - vector field

Two basic examples

- ① $L(z)$ - matrix rational functions

$$L(z) = \sum \frac{u_i}{z-z_i} + u_0 \quad u_0, u_i \in \text{Mat}_{r \times r}$$

Particular case of the Hitchin system

- Lax equation \Rightarrow integrals of motion
- Lax equation \Rightarrow "defines" A

$$\det(\kappa \cdot \text{Id} - L(z)) = \kappa^r + \sum_{i=1}^r s_i(z) \kappa^{r-i}$$

$s_i(z)$ - rational function of z

$$\dot{s}_i(z) = 0$$

$$\dots \dots \dots \dots \dots N(z) = 1$$

$$\dot{s}_i(z) = 0$$

Consider $(\partial_t - A) \Psi(z, t) = 0 \quad \Psi(z, 0) = 1$

$$(\partial_t - A) L(t) \Psi(z, t) = 0$$

$$L(t) \Psi(t) = \Psi(t) G$$

$$t=0 \Rightarrow G = L(0)$$

$$L(t) = \Psi(t) L(0) \Psi^{-1}(t)$$

{integrals}

$s_i(z)$ has poles of order i at z_i

$$N \left(\sum_{i=1}^r i \right) + r = N \frac{r(r+1)}{2} + r$$

$$\dim(L(z)) = Nr^2 + r^2$$

Lax equations are invariant under gauge transform

$$L \rightarrow g L g^{-1} \quad A = g A g^{-1} \quad g - \text{const}$$

$$\dim(L / \mathcal{C}L_r) = Nr^2 + 1$$

Consider $L(z) = \sum \frac{u_i}{z - z_i} + u_0 \quad u_0^{ij} = u_0^i \delta_{ij}$

For fixed u_0 $\dim L = Nr^2 + r$

Equations

$\dot{L} = [A, L]$ $[A, L]$ - should be a rational function with the same divisors and poles

rational function with the same divisor of poles

$$A = \frac{v}{z-\mu} \rightarrow [v, L(\mu)] = 0$$

$$\mu \neq z_i \quad v = L^n(\mu)$$

$$\partial_{t_{n,\mu}} u_i = \left[\frac{L^n(\mu)}{z_i - \mu}, u_i \right] \quad u_i, \dot{u}_0 = 0$$

$$\cdot \left[\partial_{t_{n,\mu}}, \partial_{t_{n',\mu'}} \right] = 0$$

$$[\partial_\alpha - A_\alpha, \partial_\beta - A_\beta] = 0$$

$$\alpha = (n, \mu) \\ \beta = (n', \mu')$$

$$\partial_\beta A_\alpha - \partial_\alpha A_\beta + [A_\alpha, A_\beta] = 0$$

$$\partial_{t_{n',\mu'}} L^n(\mu) + \left[L^n(\mu), \frac{L^{n'}(\mu')}{\mu - \mu'} \right] = 0$$

$$\partial_{t_{n',\mu'}} L^n(z) = \left[\frac{L^{n'}(\mu')}{z - \mu'}, L^n(z) \right] \Big|_{z=\mu}$$

- Lax equation by itself "defines" $A = A(L)$
(ambiguity in the definition leads to commuting flows)

ξ_v Rigid body

$$L = z a + U$$

$$A = z b + V$$

$$a^i = a_i \delta_{ij} \quad a^i = \underline{\underline{\text{const}}}$$

$$\dot{U} = \dot{L} = [z b + V, z a + U] = z^2 [b, a] + z \left([V, a] + [b, U] \right)$$

$$\dot{U} = \dot{L} = [z b + V, z a + U] = z^2 [b, a] + z ([V, a] + [b, U]) + [V, U]$$

$$[b, a] = 0 \quad b^{ij} = b_i \delta_{ij} \quad a_i \neq a_j$$

$$V^{ij} (a_i - a_j) = (b_i - b_j) U^{ij}$$

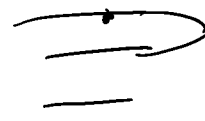
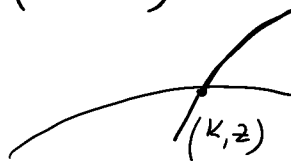
$$V^{ij} = \frac{b_i - b_j}{a_i - a_j} U^{ij} \quad b_i = a_i^2$$

$$V^{ij} = (a_i + a_j) U^{ij} \quad V = a U + U a$$

$$\dot{U} = [a U + U a, U]$$

Next step \rightarrow spectral transform line bundle
 $L(z) \rightarrow$ (spectral curve, ?) of eigenvector

$$0 = \det(\kappa - L(z)) \subset \mathbb{C}^2(\kappa, z)$$



$L(z)$

(Γ)

Spectral curves

z